# Cosmic Rays and Gamma Rays in the InterStellar Medium



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# The discovery of Cosmic Rays

At the beginning of the 20th century, the discharge rate of an electroscope was used as a measure of the level of radioactivity

electroscopes discharge slowly even in the absence of a radioactive source -> background radiation

\* radiation from radioactive materials in the Earth?



#### ELECTROSCOPE

# The discovery of Cosmic Rays

If due to radioactive materials in the Earth, the effect should diminish with height

In 1912, during a balloon flight Victor Hess discovered that the effect was indeed **increasing with height**, and concluded that:

"a radiation of very high penetrating power enters our atmosphere from above"



V. Hess in 1912

#### What are Cosmic Rays?

Cosmic rays particles hit the Earth's atmosphere at the rate of about 1000 **per square meter per second**. They are ionized nuclei - about 90% **protons**, 9% alpha particles and the rest heavy nuclei - and they are distinguished by their high energies. Most cosmic rays are **relativistic**, having energies comparable or somewhat greater than their masses. A very few of them have ultrarelativistic energies extending up to  $10^{20}$  eV (about 20 Joules), eleven order of magnitudes greater than the equivalent rest mass energy of a proton. The fundamental question of cosmic ray physics is, "Where do they come from?" and in particular, "How are they accelerated to such high energies?".

T. Gaisser "Cosmic Rays and Particle Physics"

Also electrons are present in the cosmic radiation -> ~ 1%



































Cosmic Ray anisotropy:  $\delta = \frac{I_{max} - I_{min}}{I_{max} + I_{min}}$ 

(I -> CR intensity)



figure from Iyono et al, 2005

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#### Energy density

Cosmic Ray energy density: 
$$w_{CR} \sim 1 \text{ eV cm}^{-3}$$
  
Magnetic field energy density:  $w_B = \frac{B^2}{8\pi} \sim 1 \text{ eV cm}^{-3}$   
Thermal gas energy density:  $w_{gas}^{turb} = \rho_{gas} v_{turb}^2 \sim 1 \text{ eV cm}^{-3}$ 

CRs are dynamically important in the Galaxy

#### Variations in time and space

CR flux at Earth constant during the last 10° yr (from radiation damages in geological and biological samples, meteorites, and lunar rocks)
thus the CR flux must be constant along the orbit of the Sun around the galactic centre (many revolutions in a Gyr)

Stability in time and (hints for) spatial homogeneity

#### What we have to explain about CRs:

Energy density

🧆 Energy spectrum

Chemical composition

🥯 Isotropy

🥯 Stability in time

Spatial homogeneity (?)

#### Cosmic Rays and Gamma-Ray Astronomy

Cosmic Rays undergo hadronic interactions in the InterStellar Medium:



 $\pi^0 \to \gamma + \gamma$ 

The gamma ray emission traces the gas distribution (times the CR distribution)

### Cosmic Rays and Gamma-Ray Astronomy



Energy threshold for neutral pion production:

 $p + p \rightarrow p + p + \pi^0$ 



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$$\left(E_p + m_p c^2\right)^2 - p_p^2 c^2 = E^2 - p^2 c^2 = \left(2m_p c^2 + m_{\pi^0} c^2\right)^2$$

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$$E_p - m_p c^2 > 2m_{\pi^0} c^2 + \left(\frac{m_{\pi^0}}{2m_p}\right) m_{\pi^0} c^2 \approx 280 \text{ MeV}$$

CRs produce gammas

Let's calculate the spectrum of neutral pions:

We assume a power law spectrum for CRs:  $N_p(E_p) \propto E_p^{-\delta}$ 

Fraction of proton kinetic energy transferred to pion (from data):  $f_{\pi^0}pprox 0.17$ 

production rate  $q_{\pi^0} = \int dE_p \ N_p(E_p) \ \delta(E_{\pi^0} - f_{\pi^0}E_{p,kin}) \ \sigma_{pp}(E_p) \ n_{gas} \ c$ 

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Let's now calculate the spectrum of photons from pion decay - I

The photon spectrum is the result of a "one-bodydecay" (neutral pion)



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Pion rest frame:



$$E_{\gamma}^* = \frac{m_{\pi^0}}{2}$$

Lab frame:

$$E_{\gamma} = \gamma \left( E_{\gamma}^* + v p_{\gamma}^* \cos \theta^* \right)$$

max and min energies ->  $\cos \theta^* = \pm 1$ 

$$\frac{m_{\pi^0}}{2}\sqrt{\frac{1-\beta}{1+\beta}} \le E_{\gamma} \le \frac{m_{\pi^0}}{2}\sqrt{\frac{1+\beta}{1-\beta}}$$

Let's now calculate the spectrum of photons from pion decay - II

$$E_{\gamma}^{min} = \frac{m_{\pi^{0}}}{2} \sqrt{\frac{1-\beta}{1+\beta}} \le E_{\gamma} \le \frac{m_{\pi^{0}}}{2} \sqrt{\frac{1+\beta}{1-\beta}} = E_{\gamma}^{max}$$

(1) 
$$\frac{\log E_{\gamma}^{max} + \log E_{\gamma}^{min}}{2} = \log\left(\frac{m_{\pi^0}}{2}\right)$$

in log-scale, the centre of the interval is half the pion mass

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(2) in the pion rest frame the photon distribution is isotropic

$$\frac{\mathrm{d}n_{\gamma}}{\mathrm{d}\Omega^*} = const$$

$$d\Omega^* = d\phi^* d(\cos \theta^*)$$

$$E_{\gamma} = \gamma \left( E_{\gamma}^* + v p_{\gamma}^* \cos \theta^* \right) \to dE_{\gamma} \propto d(\cos \theta^*)$$

$$\int \int \frac{dn_{\gamma}}{dE_{\gamma}} = const$$
The spectrum is flat!















Let's now calculate the spectrum of photons from pion decay - III



m extstyle > the gamma ray spectrum is symmetric (in log-log) with respect to:  $rac{m_{\pi^0}}{2} \sim 70 \,\, {
m MeV}$ 

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Detectability condition for a Cherenkov Telescope

HESS sensitivity: 
$$F_{HESS}^{min}(> 1 \text{ TeV}) \approx 10^{-12} \text{ph cm}^{-2} \text{s}^{-1} \rightarrow 2 \times 10^{-12} \text{erg/cm}^2/\text{s}$$
  
 $\uparrow$   
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which corresponds to a luminosity:  $L_{HESS}^{min}(> 1 \text{ TeV}) = F_{HESS}^{min}(> 1 \text{ TeV}) 4\pi d^2 = 2 \times 10^{32} d_{kpc}^2 \text{erg/s}$ 

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to a luminosity:  $L_{HESS}^{min}(>1~\text{TeV}) = F_{HESS}^{min}(>1~\text{TeV}) \ 4\pi d^2 = 2 \times 10^{32} d_{kpc}^2 \text{erg/s} \\ & & \text{p-p energy} \\ & & \text{fraction of } CR \text{ energy} \\ & & \text{transferred to photon} \\ & & \text{which corresponds} \\ & & \text{to a CR total energy:} \\ \end{array}$   $W_{CR} = t_{pp} \ c_{p \rightarrow \gamma}^{-1} \ L_{HESS}^{min} = \\ \end{array}$ 

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HESS sensitivity:  $F_{HESS}^{min}(> 1 \text{ TeV}) \approx 10^{-12} \text{ph cm}^{-2} \text{s}^{-1} \rightarrow 2 \times 10^{-12} \text{erg/cm}^2/\text{s}$ Depends a bit on the spectrum which corresponds to a luminosity:  $L_{HESS}^{min}(>1 \text{ TeV}) = F_{HESS}^{min}(>1 \text{ TeV}) 4\pi d^2 = 2 \times 10^{32} d_{knc}^2 \text{ erg/s}$  $t_{pp} = (n_{gas} \sigma_{pp} c k)^{-1}$ inelasticity p-p energy fraction of CR energy loss time 🔨 👘 transferred to photon which corresponds to a CR total energy:  $W_{CR} = t_{pp} c_{n \rightarrow \gamma}^{-1} L_{HESS}^{min} =$  $= \left(\frac{10^{15}}{n_{gas}} \text{ s}\right) (10) (2 \times 10^{32} d_{kpc}^2 \text{ erg/s}) = (2 \times 10^{48} n_{gas}^{-1} d_{kpc}^2 \text{ erg})$ 

#### SUMMARY:

ho the gamma ray spectrum is symmetric (in log-log) with respect to:  $rac{m_{\pi^0}}{2} \sim 70 \,\, {
m MeV}$ 

> at high energy the spectrum mimics the CR spectrum, with (roughly):  $E_\gamma pprox rac{E_{CR}}{10}$ 

detectability condition:

$$W_{CR} > 2 \times 10^{48} n_{gas}^{-1} d_{kpc}^2 \text{ erg}$$

above ~ 10 TeV

Secondary electrons and positrons:

$$p + p \to p + p + \pi^0 + \pi^+ + \pi^-$$

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 $\pi^{\pm} \to \mu^{\pm} + \nu_{\mu}(\bar{\nu}_{\mu})$  $\mu^{\pm} \to e^{\pm} + \bar{\nu}_{\mu}(\nu_{\mu}) + \nu_{e}(\bar{\nu}_{e})$ 

Secondary electrons and positrons:



Final products of proton-proton interactions are not only gamma ray photons but also neutrinos, anti-neutrinos, electrons and positrons

$$E_e \approx E_\nu \approx \frac{E_p}{20}$$

# Production of gamma rays in the Galaxy

(1) proton-proton interactions:  $(lpha=\delta)$ 

CR proton spectrum:  $N_{CR} \propto E^{-2.7}$   $\longrightarrow$  gamma ray spectrum:  $q_\gamma \propto E^{-2.7}$ 

(plus the pion bump at ~70 MeV)

Production of gamma rays in the Galaxy (1) proton-proton interactions:  $(\alpha = \delta)$ CR proton spectrum:  $N_{CR} \propto E^{-2.7}$   $\implies$  gamma ray spectrum:  $q_\gamma \propto E^{-2.7}$ (plus the pion bump at ~70 MeV) (2) inverse Compton scattering:  $(\alpha = \frac{\delta + 1}{2})$ CR electron spectrum:  $N_e \propto E^{-3}$  gamma ray spectrum:  $q_\gamma \propto E^{-2}$ 

(plus the Klein-Nishina cutoff at ~50 TeV)

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CR electron spectrum:  $N_e \propto E^{-3}$  rightarrow gamma ray spectrum:  $q_\gamma \propto E^{-3}$ 

# The gamma ray emission from the Galaxy

FERMI observation of the galactic diffuse emission



#### p-p interactions dominate the diffuse emission

# The gamma ray emission from the Galaxy

FERMI observation of the galactic diffuse emission



CRs quite homogeneous in the Galaxy!

#### Take-home message

We have plenty of data on CRs but we still don't

know where they are from;

A connection exists between CR physics and Gamma
Ray Astronomy because CRs produce gamma rays in
interactions with matter and radiation fields;
What we know about CRs seems to explain fairly
well the gamma ray diffuse emission we observe from

the Milky Way.